Calculation of helium nuclei in quenched lattice QCD (multi-baryon systems in lattice QCD)

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1. Introduction

Spectrum of Nuclei

success of Shell model since <u>1949</u>: Jensen and Mayer degrees of freedom of protons and neutrons

Spectrum of nucleon (proton and neutron) degrees of freedom of quarks and gluons success of non-perturbative calculation of QCD such as lattice QCD

Motivation :

Understand property and structure of nuclei from QCD

If we can study nuclei from QCD, we may be able to

- 1. reproduce spectrum of nuclei
- 2. predict property of nuclei hard to calculate or observe such as neutron rich nuclei

multi-nucleon state from lattice QCD

1. NN state ${}^{3}S_{1}$ and ${}^{1}S_{0}$ '95 Fukugita et al : Quenched QCD scattering energy $\rightarrow a_0 > 0$ '06 NPLQCD : $N_f = 2 + 1$ partially quenched mixed action QCD scattering energy $\rightarrow a_0 \lesssim 0$ '07 Ishii, Aoki, Hatsuda : Quenched and $N_f = 2 + 1$ QCD wave function $\rightarrow a_0 > 0$ '09 NPLQCD : $N_f = 2 + 1$ QCD (anisotropic) scattering energy $\rightarrow a_0 \lesssim 0$ Deuteron: unbound due to $m_{\pi} \gtrsim 0.3$ GeV 2. NNN state '09 NPLQCD : $N_f = 2 + 1$ QCD (anisotropic) $\equiv^0 \equiv^0 n$ and pnn states Triton: further study necessary

Helium nucleus: larger binding energy He⁴ : double magic numbers Z = 2, N = 2

In this work : Exploratory study for He and ³He nuclei

Outline

- 1. Introduction
- 2. Problems of multi-nucleon bound state
- 3. Simulation parameters
- 4. Results for He and 3 He
- 5. Summary

2. Problems of multi-nucleon bound state

1. Statistical error

 $m_{\pi} \rightarrow \text{small}$ # nucleon \rightarrow large $\Rightarrow \frac{\text{noise}}{\text{signal}} \rightarrow$ large Avoid large statistical fluctuation unphysically heavy quark mass + $O(10^3)$ measurements

 $m_{\pi} = 0.8$ GeV and $m_N = 1.62$ GeV

2. Calculation cost

3. Identification of bound state in finite volume

Calculation cost

 $C_{\text{He}}(t) = \langle 0 | \text{He}(t) \overline{\text{He}}(0) | 0 \rangle$ with $\text{He} = p^2 n^2 = [udu]^2 [dud]^2$

Number of Wick contraction $N_u! \times N_d! = (2N_p + N_n)! \times (2N_n + N_p)!$ contain identical contractions

> He: $6! \times 6! = 518400$ ³He: $5! \times 4! = 2880$

Reduction of contractions

Symmetries $p \leftrightarrow p, n \leftrightarrow n$ in He operator Isospin all $p \leftrightarrow$ all nCalculate two contractions simultaneously

 $u \leftrightarrow u$ in p or $d \leftrightarrow d$ in n

Calculation cost (cont'd)

 $C_{\text{He}}(t) = \langle 0 | \text{He}(t) \overline{\text{He}}(0) | 0 \rangle$ with $\text{He} = p^2 n^2 = [udu]^2 [dud]^2$

Number of Wick contraction $N_u! \times N_d! = (2N_p + N_n)! \times (2N_n + N_p)!$ contain identical contractions

He:
$$6! \times 6! = 518400 \longrightarrow 1107$$

³He: $5! \times 4! = 2880 \longrightarrow 93$

Furthermore, avoid same calculations of dirac and color indices

Block of three quark propagators B_3

zero momentum nucleon operator in sink time slice

Blocks of two B_3

1, 2, 3 dirac contractions carried out

Identification of bound state in finite volume

Example) Two-particle system

observe small $\Delta E = E - 2m < 0$ at single L



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Identification of bound state in finite volume (cont'd) Example) Two-particle system

observe small $\Delta E = E - 2m < 0$ at single L



Bound state : $\Delta E = -\Delta E_{bind} + O(e^{-\gamma L}) < 0$ Beane *et al.*, PLB585:106(2004), Sasaki, TY, PRD74:114507(2006)

Identification of bound state in finite volume (cont'd) Example) Two-particle system



Identification of bound state in finite volume (cont'd) Example) Two-particle system

observe small
$$\Delta E = E - 2m < 0$$
 at several L



Identify bound state from volume dependence of ΔE observe constant in infinite volume limit with L = 3.1, 6.1, 12.3 fm

3. Simulation parameters

- Quenched Iwasaki gauge action at $\beta = 2.416$ $a^{-1} = 1.54~{\rm GeV}$ with $r_0 = 0.49~{\rm fm}$
- Tad-pole improved Wilson fermion action

$$m_{\pi} = 0.8$$
 GeV and $m_N = 1.62$ GeV

• Three volumes

L	L [fm]	N _{conf}	N _{meas}
24	3.1	2500	2
48	6.1	400	12
96	12.3	200	12

• Exponential smearing sources $q(\vec{x}) = A \exp(-B|\vec{x}|)$

$$S_1$$
 S_2
(A, B) = (0.5, 0.5), (0.5, 0.1) for $L = 24$
(A, B) = (0.5, 0.5), (1.0, 0.4) for $L = 48,96$

• quark operator with non-relativistic projection in nucleon operator

Simulations:

PACS-CS at Univ. of Tsukuba, and HA8000 at Univ. of Tokyo

4. Results

Effective mass of He and ³He nuclei at L = 48 $m_{\text{He}}(t) = \log \left(\frac{C_{\text{He}}(t)}{C_{\text{He}}(t+1)}\right)$



- Clear signal in t < 12 and larger error in $t \ge 12$
- consistent plateaus in $8 \lesssim t \le 12$

4. Results (cont'd)

Effective energy shift $\Delta E_L = m_{\text{He},^3\text{He}} - Nm_N$ of He nuclei at L = 48 $\Delta E_L(t) = \log\left(\frac{R(t)}{R(t+1)}\right), \quad R_{\text{He}}(t) = \frac{C_{\text{He}}(t)}{(C_N(t))^4}, \quad R_{^3\text{He}}(t) = \frac{C_{^3\text{He}}(t)}{(C_N(t))^3}$



•
$$\Delta E_L < 0$$
 in $8 \lesssim t \le 12$

• consistent plateaus in $8 \lesssim t \le 12$

4. Results (cont'd)

Volume dependence of $\Delta E_L = -\Delta E_{bind} + F(L)$ of He nuclei



- Small volume dependence
- Infinite volume limit with $F(L) = C/L^3$
- Non-zero binding energy in infinite volume limit

4. Results (cont'd)

Volume dependence of $\Delta E_L = -\Delta E_{bind} + F(L)$ of He nuclei



Binding increases as mass number in experiment, but inconsistent $\Delta E_{\text{He}}/4 = 6.9(2.0)(1.4) \text{ MeV}$ and $\Delta E_{^{3}\text{He}}/3 = 6.1(1.2)(1.0) \text{ MeV}$ mainly caused by heavy quark mass in calculation, probably

5. Summary

- Exploratory study of helium nuclei in quenched lattice QCD
- Unphysically heavy quark mass
- Reduction of calculation cost with some techniques
- Volume dependence of energy shift from free multi-nucleon state

Non-zero energy shift in infinite volume limit \rightarrow He and ³He are bound at $m_{\pi} = 0.8$ GeV

Future work

- Quark mass dependence of ΔE
- Reduction of statistical error
- Deuteron bound state
- Larger nuclei
- Dynamical quark effect